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I. Magnetic symmetry

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BILBAO, SPAIN

Symmetry-Based Computational Tools for Magnetic Crystallography

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[Annu. Rev. Mater. Res. 2015. 45:217–48](#)

[DOI: 10.1146/annurev-matsci-070214-021008](#)

Magnetic Section



bilbao crystallographic server

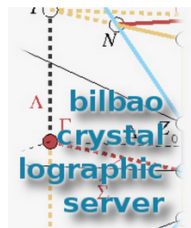


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About us

Publications

How to cite the server



Crystallography Online: Workshop on the use of the structural and magnetic tools of the Bilbao Crystallographic Server
September 2021, Leioa (Spain)

Forthcoming schools and workshops

News:

- **New Article in Nature**
10/2020: Xu *et al.* "High-throughput calculations of magnetic topological materials" *Nature* (2020) **586**, 702-707.
- New programs: **MBANDREP**, **COREPRESENTATIONS**, **COREPRESENTATIONS PG**, **MCOMPAREL**, **MSITESYM**, **MKVEC**, **Check Topological Magnetic Mat**
10/2020: new tools in the sections "Magnetic Symmetry and Applications" and "Representations and Applications". [More info](#)

Space-group symmetry

Magnetic Symmetry and Applications

Group-Subgroup Relations of Space Groups

Representations and Applications

Solid State Theory Applications

Structure Utilities

Topological Quantum Chemistry

Subperiodic Groups: Layer, Rod and Frieze Groups

Structure Databases

Raman and Hyper-Raman scattering

Quick access to some tables

Space Groups

Plane Groups

Layer Groups

Rod Groups

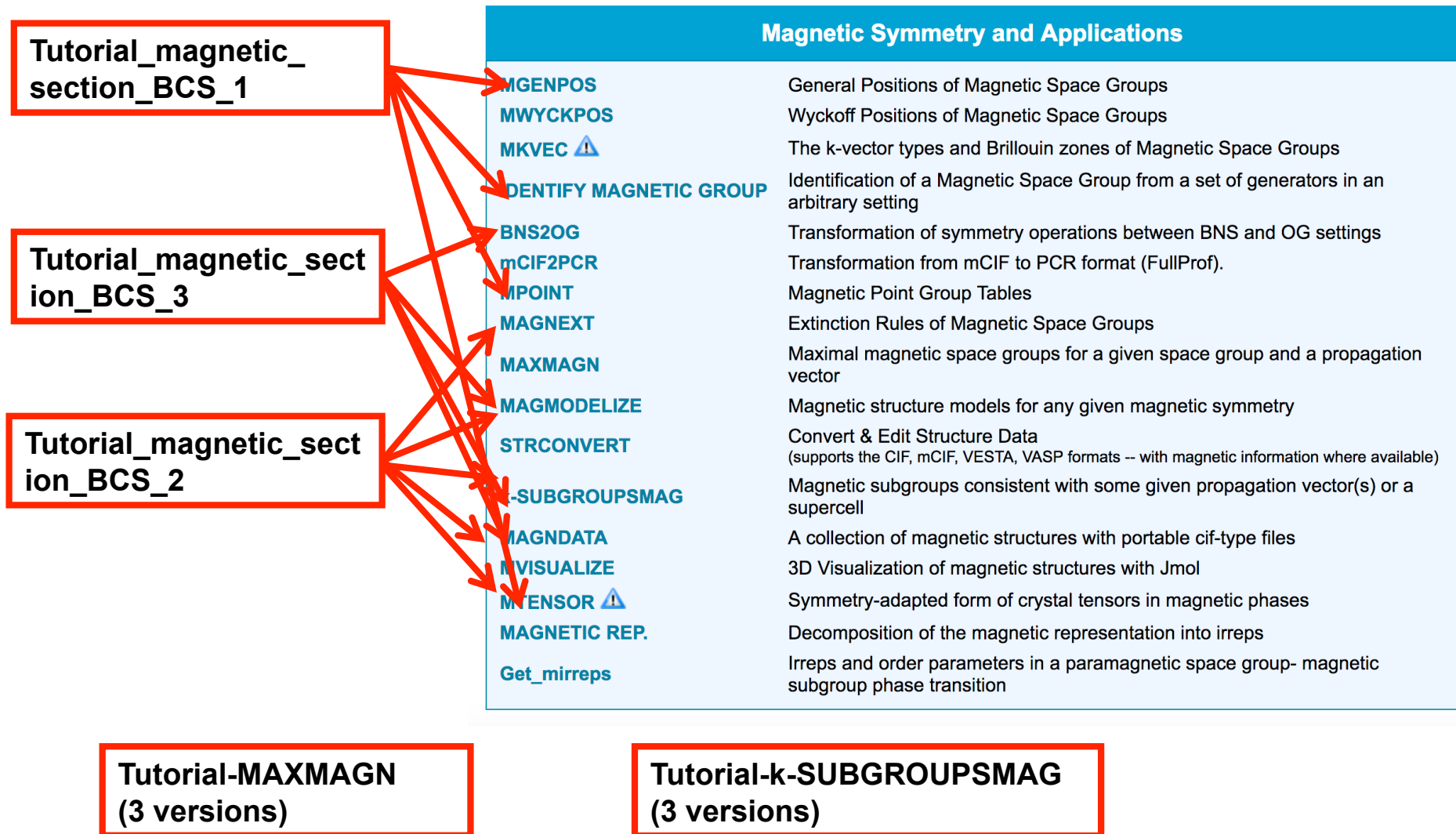
Frieze Groups

2D Point Groups

3D Point Groups

Magnetic Space Groups

Three main tutorials on the programs of the BCS Magnetic Section can be directly downloaded from the webpages of the programs :

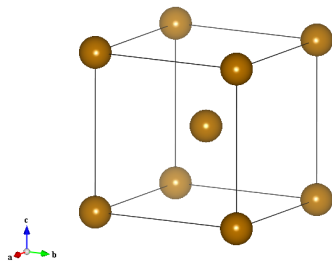


Symmetry and Physics:

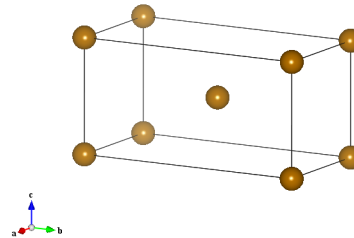
A symmetry property in a solid is **NOT ONLY** some mathematical property. It is a **PHYSICAL PROPERTY!**

A well defined symmetry operation in a thermodynamic system must be maintained when scalar fields (temperature, pressure,...) are changed, **except if a phase transition takes place.**

The change of symmetry of a crystal necessarily implies a phase transition.



$a=b=c$ symmetry property



$a = c$
 $b = 2a$

"nice" but not a symmetry property

Defining the symmetry of a crystal

1st step. We define a set of operations/transformations on the system: rotations, translations, space inversion, ETC. (they form a mathematical group)

2nd step. On the previous group of operations we look for the subset (subgroup) of operations, which keep the crystal INVARIANT or UNDISTINGUISHABLE . This is the symmetry group of the crystal.

But what is the group of operations that we define or choose in the first step above? (Once this group is chosen, there is no ambiguity on the symmetry group of the system). **here comes the Physics:**

Symmetry and Physics:

Group of all possible combinations of **rotations**
translations
space inversion

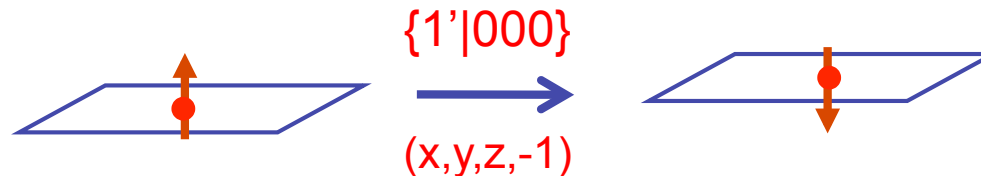
(They all keep the **ENERGY of the system invariant....**)

The symmetry group of the solid is formed by all operations **that keeping the ENERGY invariant ALSO** maintain the system undistinguishable after applying the operation.

The time reversal operation also keeps energy invariant:

Definition of time reversal: $\{1'|0,0,0\}$:

- Does not change nuclear variables
- Changes sign of ALL atomic (average) magnetic moments



If all average atomic moments are zero, the system is invariant for the time reversal operation:

Time reversal symmetry is present as symmetry operations in non-magnetic structures but it is ABSENT in magnetically ordered ones!

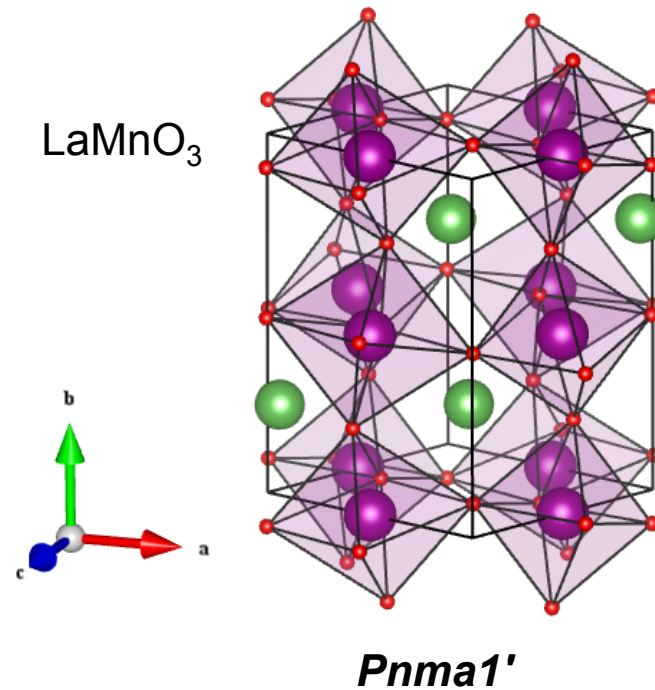
Magnetic symmetry groups:

We do not add but SUBTRACT symmetry operations !

A symmetry operation is detected when it does NOT exist !

All NON-magnetic structures have time reversal symmetry!

All symmetry operations are (implicitly) duplicated: with and without time reversal



$(x, y, z, +1)$	$(-x+1/2, -y, z+1/2, +1)$	$(-x, y+1/2, -z, +1)$	$(x+1/2, -y+1/2, -z+1/2, +1)$
$(-x, -y, -z, +1)$	$(x+1/2, y, -z+1/2, +1)$	$(x, -y+1/2, z, +1)$	$(-x+1/2, y+1/2, z+1/2, +1)$
$(x, y, z, -1)$	$(-x+1/2, -y, z+1/2, -1)$	$(-x, y+1/2, -z, -1)$	$(x+1/2, -y+1/2, -z+1/2, -1)$
$(-x, -y, -z, -1)$	$(x+1/2, y, -z+1/2, -1)$	$(x, -y+1/2, z, -1)$	$(-x+1/2, y+1/2, z+1/2, -1)$

All NON-magnetic structures have time reversal symmetry

If all atomic magnetic moments are zero, time inversion is a (trivial) symmetry operation of the structure:

Actual symmetry of the non-magnetic phase:

$$Pnma1' = Pnma + \{1'|000\} \times Pnma \quad (\text{grey group})$$

16 operations:

$$\begin{array}{cccc}
 (x,y,z,+1) & (-x+1/2,-y,z+1/2,+1) & (-x,y+1/2,-z,+1) & (x+1/2,-y+1/2,-z+1/2,+1) \\
 (-x,-y,-z,+1) & (x+1/2,y,-z+1/2,+1) & (x,-y+1/2,z,+1) & (-x+1/2,y+1/2,z+1/2,+1) \\
 (x,y,z,-1) & (-x+1/2,-y,z+1/2,-1) & (-x,y+1/2,-z,-1) & (x+1/2,-y+1/2,-z+1/2,-1) \\
 (-x,-y,-z,-1) & (x+1/2,y,-z+1/2,-1) & (x,-y+1/2,z,-1) & (-x+1/2,y+1/2,z+1/2,-1)
 \end{array}$$

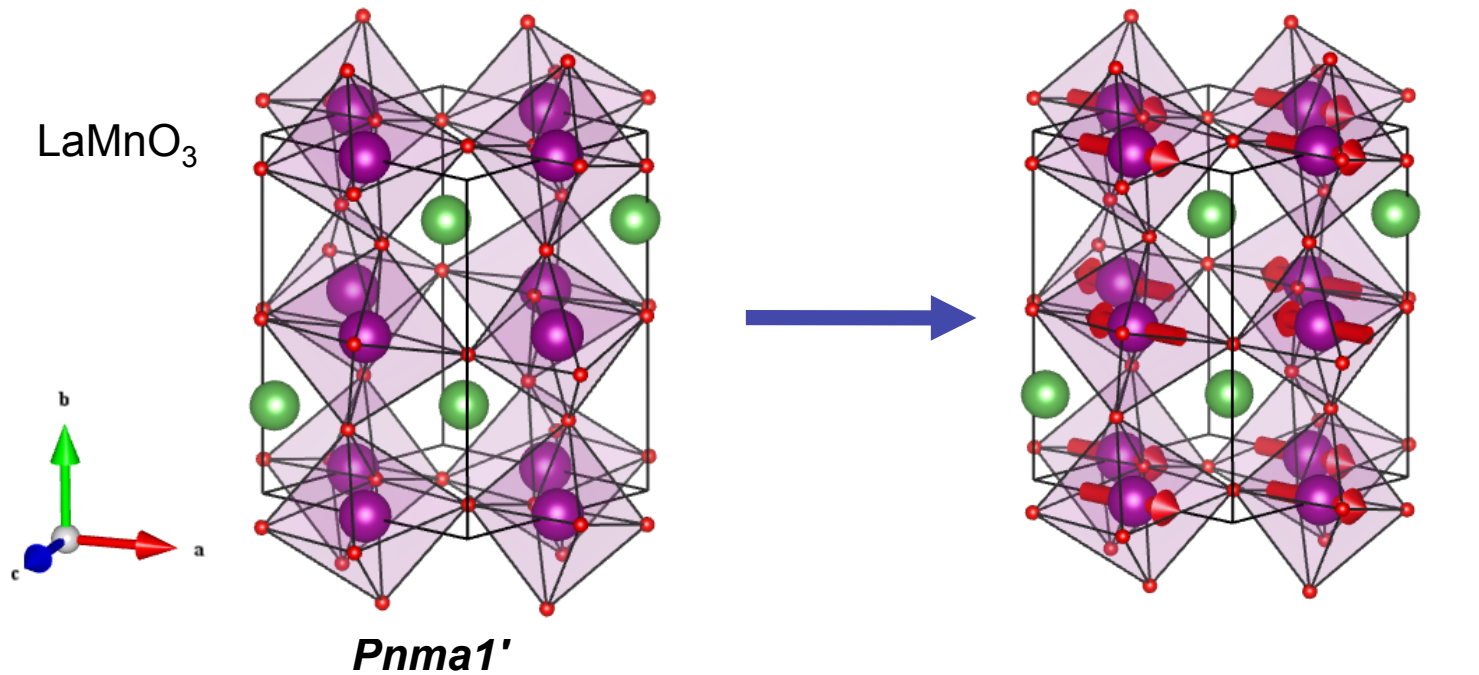
Notation:

$$\begin{array}{l}
 (x+1/2,-y+1/2,-z+1/2,+1) == \{2x| \frac{1}{2} \frac{1}{2} \frac{1}{2} \} \{R|t\} \\
 (x+1/2,-y+1/2,-z+1/2,-1) == \{2x| \frac{1}{2} \frac{1}{2} \frac{1}{2} \} \{R'|t\}
 \end{array}
 \quad \{R,\theta|t\} \begin{array}{l} \rightarrow \theta=1 \\ \rightarrow \theta=-1 \end{array}$$

Magnetic ordering is a symmetry breaking process

Magnetically ordered phases:

Time reversal $\{1'|0\ 0\ 0\}$ is LOST

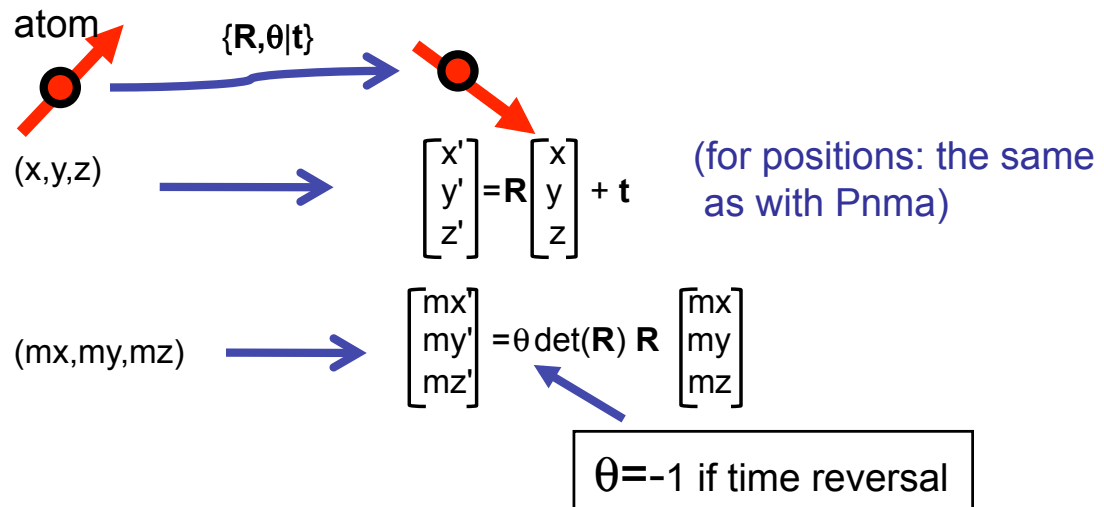
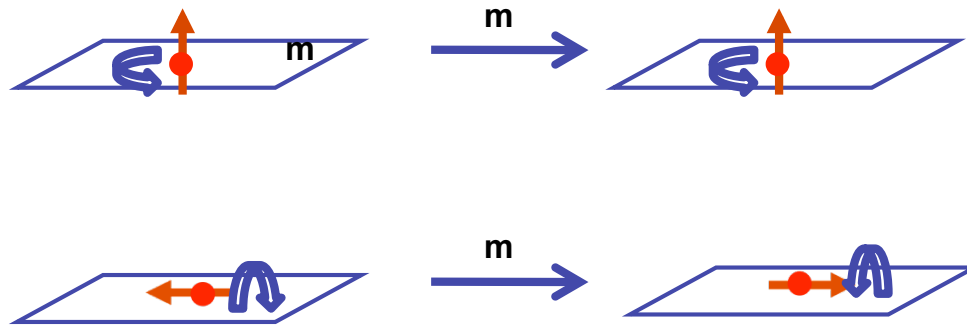


$(x,y,z,+1)$	$(-x+1/2,-y,z+1/2)$	$(x+1/2,-y+1/2,-z+1/2,+1)$
$(-x,-y,-z,+1)$	$(x+1/2,y,-z+1/2)$	$(x+1/2,y+1/2,z+1/2,+1)$
$(x,y,z,-1)$	$(-x+1/2,-y,z+1/2)$	$(x+1/2,-y+1/2,-z+1/2,-1)$
$(-x,-y,-z,-1)$	$(x+1/2,y,-z+1/2)$	$(x+1/2,y+1/2,z+1/2,-1)$

?

For space operations, the magnetic moments transform as pseudovectors or axial vectors:

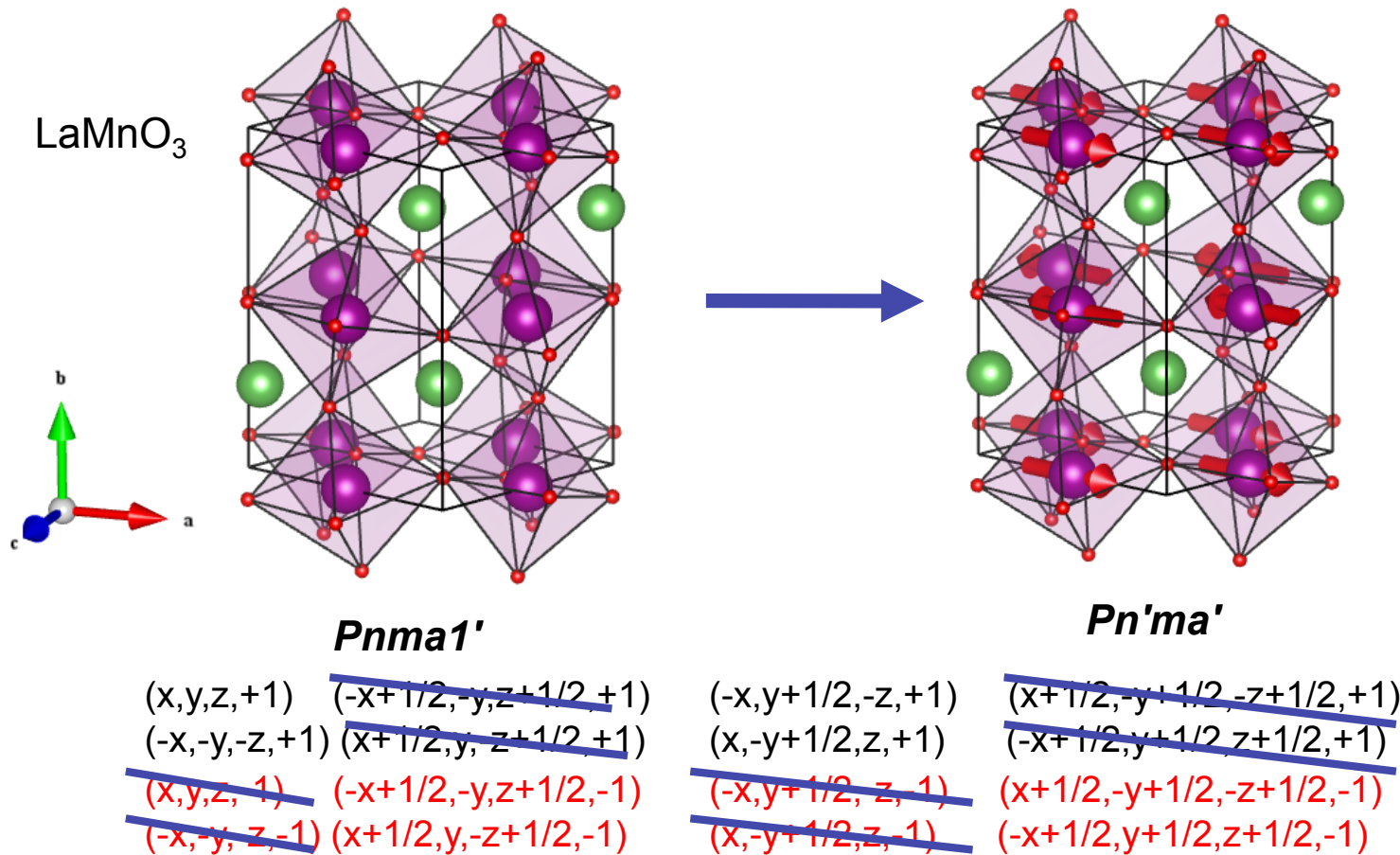
$$T_{\text{axial}}(\mathbf{R}) = \det[\mathbf{R}] \mathbf{R}$$



Magnetic ordering is a symmetry breaking process

Magnetically ordered phases:

Time reversal $\{1'|0\ 0\ 0\}$ is LOST
but operations including time reversal can be maintained

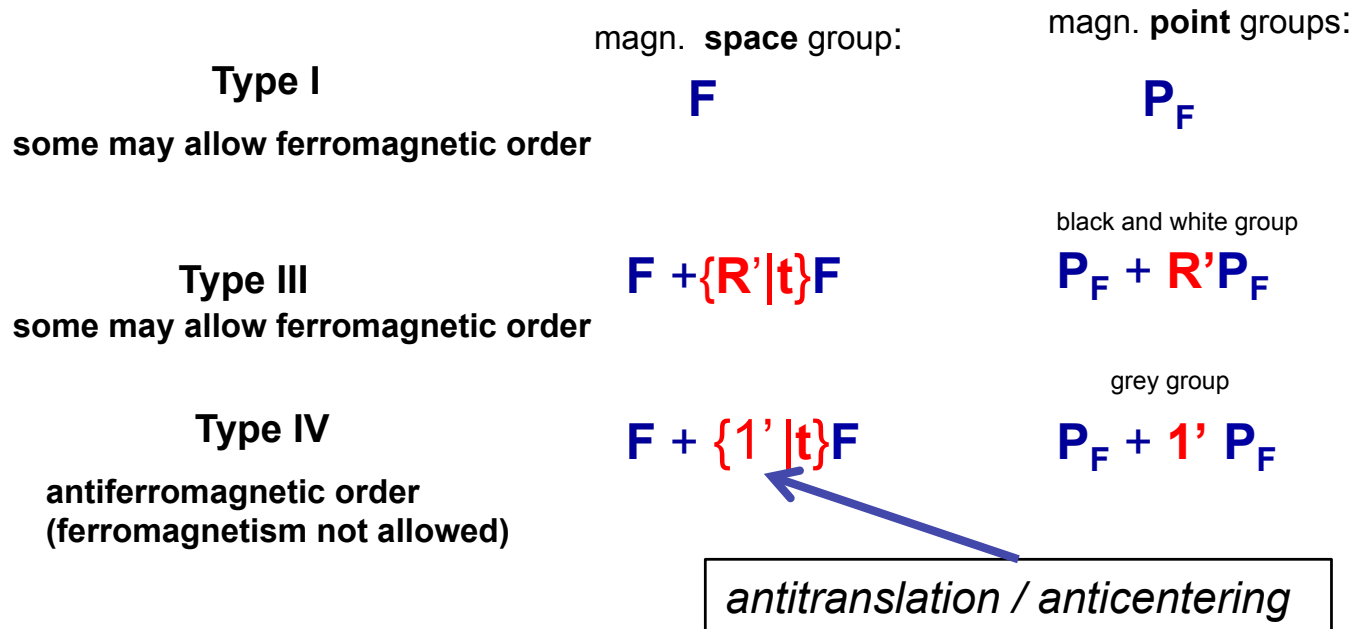


Types of magnetic space groups:

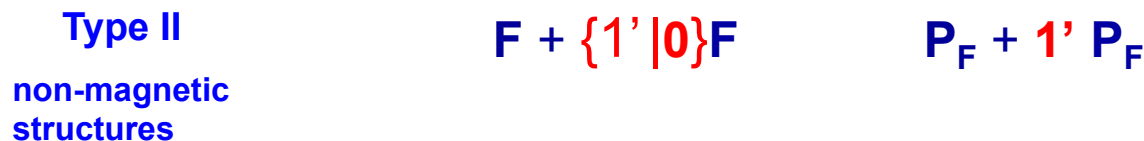
(for a commensurate magnetic structure resulting from a paramagnetic phase having a grey magnetic group $G1'$)

F subgroup of G
 $F \leq G$

Time reversal $\{1' | 0 0 0\}$ is NOT a symmetry operation of a magnetic structure, but combined with a translation it can be...



(Type II are the grey groups of the non-magnetically ordered systems):



Tables of magnetic space groups ("standard" settings)

1.- e-book: D.B. Litvin: "Magnetic space groups" (Electronic Book)

Litvin DB. 2013. *Magnetic Group Tables: 1-, 2- and 3-Dimensional Magnetic Subperiodic Groups and Magnetic Space Groups*. Chester, UK: Int. Union Crystallogr. <http://www.iucr.org/publ/978-0-9553602-2-0>

(listing using only OG setting)

2.- Computer readable listing:

ISOTROPY webpage: <http://stokes.byu.edu/iso/magneticspacegroups.html>

H.T. Stokes and B.J. Campbell

(downloadable files using BNS and OG settings)

3.- Web online listing: Bilbao crystallographic server (www.cryst.ehu.es)

Magnetic Symmetry and Applications

MGENPOS

General Positions of Magnetic Space Groups

MWYCKPOS

Wyckoff Positions of Magnetic Space Groups

(listings using
BNS and OG settings)

(So far) only software using BNS setting exists

Fundamental difference of the OG description:

For type IV MSGs it uses a unit cell which does NOT describe the actual lattice of the system.

NEW MSG SYMBOLS



FOUNDATIONS
ADVANCES

ISSN 2053-2733

Introducing a unified magnetic space-group symbol

Branton J. Campbell,^{a*} Harold T. Stokes,^a J. Manuel Perez-Mato^b and Juan Rodríguez-Carvajal^c

Acta Cryst. (2022). *A78*, 99–106

“...., a new unified (UNI) MSG symbol is introduced, which combines a modified BNS symbol with essential information from the OG symbol.”

BNS No.	BNS	OG No.	OG	UNI MSG	UNI MPG
2.7	$P_5\bar{1}$	2.4.7	$P_{2s}\bar{1}$	$P\bar{1}.1'_c[P\bar{1}]$	$\bar{1}.1'$
42.223	F_5mm2	25.9.163	P_1mm2	$Fmm2.1'_l[Pmm2]$	$mm2.1'$
161.72	R_13c	160.5.1299	R_R3m'	$R3c.1'_c[R3m]$	$3m.1'$
218.84	$P_1\bar{4}3n$	217.5.1584	$I_p\bar{4}'3m'$	$P\bar{4}3n.1'_l[I\bar{4}3m]$	$\bar{4}3m.1'$
140.550	I_c4/mcm	123.19.1017	$P_14/m'm'm'$	$I4/mcm.1'_c[rP4/mmm]$	$4/mmm.1'$
28.96	P_Bma2	39.7.284	A_pbm2	$Pma2.1'_B[Bma2]$	$mm2.1'$

General Positions of the Group $Pn'ma'$ (#62.448)

For this space group, BNS and OG settings coincide.
Its label in the OG setting is given as: $Pn'ma'$ (#62.8.509)

N	Standard/Default Setting			
	(x,y,z) form	Matrix form	Geom. interp.	Seitz notation
1	x, y, z, +1 m_x, m_y, m_z	$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}$	$1 \underline{+1}$	{ 1 0 }
2	-x, y+1/2, -z, +1 $-m_x, m_y, -m_z$	$\begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 1/2 \\ 0 & 0 & -1 & 0 \end{pmatrix}$	2 (0, 1/2, 0) 0, y, 0 $\underline{+1}$	{ 2_{010} 0 1/2 0 }
3	-x, -y, -z, +1 m_x, m_y, m_z	$\begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \end{pmatrix}$	-1 0, 0, 0 $\underline{+1}$	{ -1 0 }
4	x, -y+1/2, z, +1 $-m_x, m_y, -m_z$	$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 1/2 \\ 0 & 0 & 1 & 0 \end{pmatrix}$	m x, 1/4, z $\underline{+1}$	{ m_{010} 0 1/2 0 }
5	$x+1/2, -y+1/2, -z+1/2, -1$ $-m_x, m_y, m_z$	$\begin{pmatrix} 1 & 0 & 0 & 1/2 \\ 0 & -1 & 0 & 1/2 \\ 0 & 0 & -1 & 1/2 \end{pmatrix}$	2 (1/2, 0, 0) x, 1/4, 1/4 $\underline{-1}$	{ $2'_{100}$ 1/2 1/2 1/2 }
6	$-x+1/2, -y, z+1/2, -1$ $m_x, m_y, -m_z$	$\begin{pmatrix} -1 & 0 & 0 & 1/2 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1/2 \end{pmatrix}$	2 (0, 0, 1/2) 1/4, 0, z $\underline{-1}$	{ $2'_{001}$ 1/2 0 1/2 }
7	$-x+1/2, y+1/2, z+1/2, -1$ $-m_x, m_y, m_z$	$\begin{pmatrix} -1 & 0 & 0 & 1/2 \\ 0 & 1 & 0 & 1/2 \\ 0 & 0 & 1 & 1/2 \end{pmatrix}$	n (0, 1/2, 1/2) 1/4, y, z $\underline{-1}$	{ m'_{100} 1/2 1/2 1/2 }
8	$x+1/2, y, -z+1/2, -1$ $m_x, m_y, -m_z$	$\begin{pmatrix} 1 & 0 & 0 & 1/2 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 1/2 \end{pmatrix}$	a x, y, 1/4 $\underline{-1}$	{ m'_{001} 1/2 0 1/2 }

Go to the list of the Wyckoff Positions of the Group $Pn'ma'$ (#62.448)
Go to the Systematic Absences for the Group $Pn'ma'$ (#62.448)

Output of
MGENPOS

Example
of type III
MSG

Magnetic
point group:
 $m'mm'$

$$Pn'ma' = P12_1/m1 + \{2'_{100}|1/2,1/2,1/2\} P12_1/m1$$

General Positions of the Group P_bmn2_1 (#31.129) [BNS setting]

To display the general positions in the OG setting, please follow this link: [P_{2b}mn2₁ \(#31.6.217\) \[Transformation matrix\]](#)

Translation lattice generators: (1|1,0,0), (1|0,1,0), (1|0,0,1), (1|0,0,0)

Black-and-white lattice generators: (1|1,0,0), (1|0,1,0), (1|0,0,1), (1'|0,1/2,0)

N	Standard/Default Setting			
	(x,y,z) form	Matrix form	Geom. interp.	Seitz notation
1	x, y, z, +1 m _x ,m _y ,m _z	$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}$	1 <u>+1</u>	{ 1 0 }
2	-x+1/2, -y, z+1/2, +1 -m _x , -m _y , m _z	$\begin{pmatrix} -1 & 0 & 0 & 1/2 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1/2 \end{pmatrix}$	2 (0,0,1/2) 1/4,0,z <u>+1</u>	{ 2 ₀₀₁ 1/2 0 1/2 }
3	-x, y, z, +1 m _x , -m _y , -m _z	$\begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}$	m 0,y,z <u>+1</u>	{ m ₁₀₀ 0 }
4	x+1/2, -y, z+1/2, +1 -m _x , m _y , -m _z	$\begin{pmatrix} 1 & 0 & 0 & 1/2 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1/2 \end{pmatrix}$	n (1/2,0,1/2) x,0,z <u>+1</u>	{ m ₀₁₀ 1/2 0 1/2 }
5	x, y+1/2, z, -1 -m _x , -m _y , -m _z	$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 1/2 \\ 0 & 0 & 1 & 0 \end{pmatrix}$ *	t (0,1/2,0) <u>-1</u>	{ 1' 0 1/2 0 }
6	-x+1/2, -y+1/2, z+1/2, -1 m _x , m _y , -m _z	$\begin{pmatrix} -1 & 0 & 0 & 1/2 \\ 0 & -1 & 0 & 1/2 \\ 0 & 0 & 1 & 1/2 \end{pmatrix}$ *	2 (0,0,1/2) 1/4,1/4,z <u>-1</u>	{ 2' ₀₀₁ 1/2 1/2 1/2 }
7	-x, y+1/2, z, -1 -m _x , m _y , m _z	$\begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 1/2 \\ 0 & 0 & 1 & 0 \end{pmatrix}$ *	b 0,y,z <u>-1</u>	{ m' ₁₀₀ 0 1/2 0 }
8	x+1/2, -y+1/2, z+1/2, -1 m _x , -m _y , m _z	$\begin{pmatrix} 1 & 0 & 0 & 1/2 \\ 0 & -1 & 0 & 1/2 \\ 0 & 0 & 1 & 1/2 \end{pmatrix}$ *	n (1/2,0,1/2) x,1/4,z <u>-1</u>	{ m' ₀₁₀ 1/2 1/2 1/2 }

Output of
MGENPOS

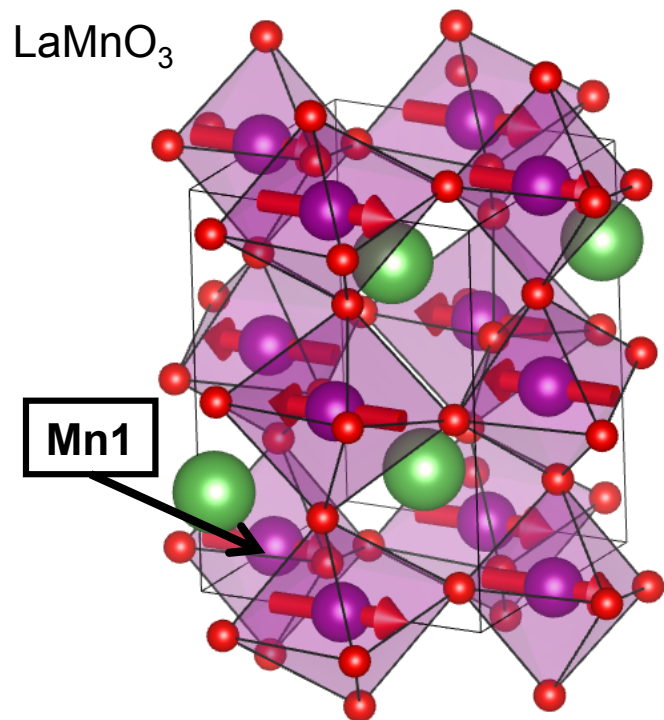
Example
of type IV
MSG

Propagation
vector $k \neq 0$

$$P_bmn2_1 = Pmn2_1 + \{1' | 0, 1/2, 0\} Pmn2_1$$

New UNI symbol:
Pmn2₁.1'_b

Description of a magnetic structure in a crystallographic form using its MSG:



Magnetic space Group:
Pn'ma'

Lattice parameters:
5.7461 7.6637 5.5333 90.000 90.000 90.000

Atomic positions of asymmetric unit:

La1	0.05130	0.25000	-0.00950	
Mn1	0.00000	0.00000	0.50000	← special position: coordinates are symmetry- forced
O1	0.48490	0.25000	0.07770	
O2	0.30850	0.04080	0.72270	

Magnetic moments of the asymmetric unit (μB) and symmetry constraints::

Mn1 3.87 0.0 0.0 (mx,my,mz)

special position BUT zero components are NOT symmetry- forced

Symmetry operations are relevant both for positions and moments

Pn' ma' :

- 1 x,y,z,+1
- 2 -x,y+1/2,-z,+1
- 3 -x,-y,-z,+1
- 4 x,-y+1/2,z,+1
- 5 x+1/2,-y+1/2,-z+1/2,-1
- 6 -x+1/2,-y,z+1/2,-1
- 7 -x+1/2,y+1/2,z+1/2,-1
- 8 x+1/2,y,-z+1/2,-1

MSG

magCIF Format

Official extension of the CIF format to communicate magnetic structures

(developed by the Commission on magnetic structures of the IUCr)

These files permit the different alternative models to be analyzed, refined, shown graphically, transported to ab-initio codes etc., with various programs as **ISODISTORT**, **JANA2006**, **STRCONVERT**, **FullProf**, **VESTA**, **Jmol**, etc.

It includes incommensurate structures !

```
-----  
_space_group_magn.transform_BNS_Pp_abc ' -b, a, c; 1/8, 1/4, 0'
```

```
-----  
_space_group_magn.number_BNS 31.129  
_space_group_magn.name_BNS "P_b m n 2_1"  
_space_group_magn.point_group_name "mm21"  
_space_group_magn.point_group_number "7.2.21"  
_cell_length_a 11.67080  
_cell_length_b 7.36060  
_cell_length_c 5.25720  
_cell_angle_alpha 90.00  
_cell_angle_beta 90.00  
_cell_angle_gamma 90.00
```

unit cell

```
loop_  
_space_group_symop_magn.operation.id  
_space_group_symop_magn.operation.xyz  
1 x, y, z, +1  
2 -x+1/4, -y, z+1/2, +1  
3 x, -y+1/2, z, +1  
4 -x+1/4, y+1/2, z+1/2, +1
```

MSG

```
loop_  
_space_group_symop_magn.centering.id  
_space_group_symop_magn.centering.xyz  
1 x, y, z, +1  
2 x+1/2, y, z, -1
```

asymmetric unit

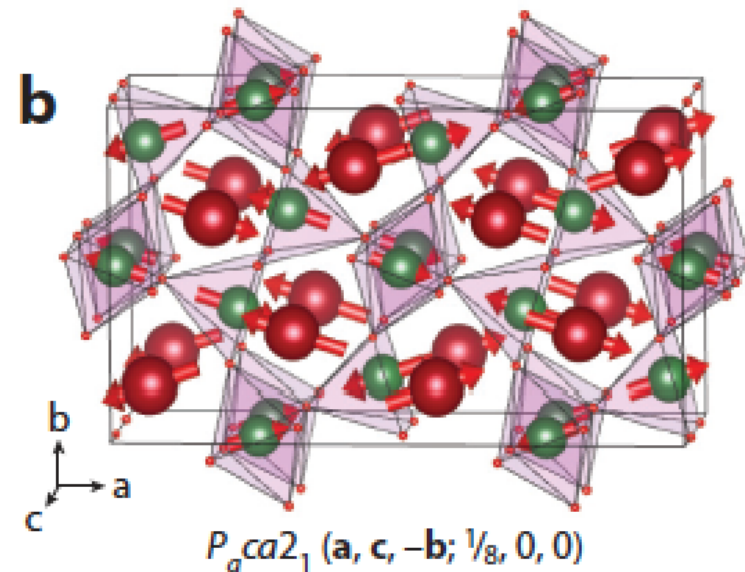
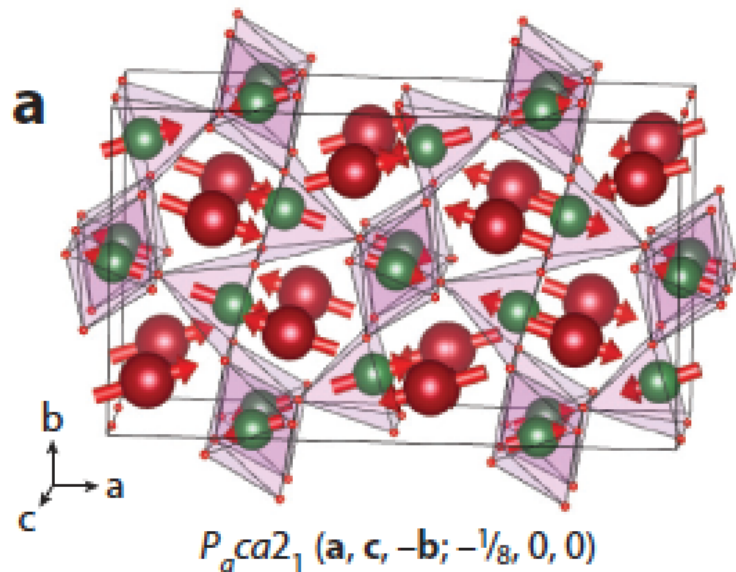
```
loop_  
_atom_site_label  
_atom_site_type_symbol  
_atom_site_fract_x  
_atom_site_fract_y  
_atom_site_fract_z  
_atom_site_occupancy  
Ho1_1 Ho 0.04195 0.25000 0.98250 1  
Ho1_2 Ho 0.95805 0.75000 0.01750 1  
Mn1 Mn1 0.00000 0.00000 0.50000 1  
O1_1 O 0.23110 0.25000 0.11130 1  
O1_2 O 0.7689 0.75000 0.88870 1  
O2_1 O 0.16405 0.05340 0.70130 1  
O2_2 O 0.83595 0.55340 0.29870 1
```

magnetic moments in the asymmetric unit

```
loop_  
_atom_site_moment.label  
_atom_site_moment.crystalaxis_x  
_atom_site_moment.crystalaxis_y  
_atom_site_moment.crystalaxis_z  
_atom_site_moment.symmform  
Mn1 3.87 0.0 0.0 mx, my, mz
```

The knowledge of the MSG allows the systematic enumeration and description of all domain-related configurations:

Gd_2MnO_5 (magndata 1.54)



twin-related spin arrangements related by space inversion, with opposite induced electric polarization.

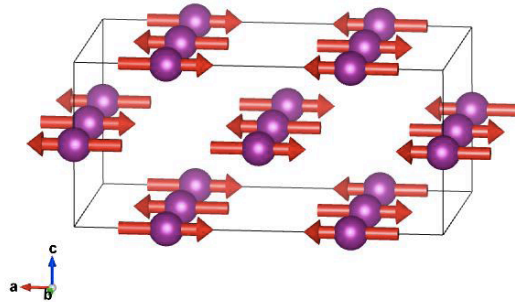
Consequences of symmetry

Von Neumann principle:

- all variables/parameters/degrees of freedom compatible with the symmetry will be present (their magnitude may be small or large, but not necessarily zero).
- Tensor crystal properties are constrained by the (magnetic) point group symmetry of the crystal.
- Reversely: any tensor property allowed by the (magnetic) point group symmetry can exist (large or small, but it is not forced to be zero)

HoMnO₃

unit cell: $2a_p, b_p, c_p$



Equivalent to the use of space group **Pnm2₁(31)** with **half cell along a**:

Atomic positions of asymmetric unit:

Ho1	4a	0.04195	0.25000	0.98250
Ho2	4a	0.95805	0.75000	0.01750
Mn1	8b	0.00000	0.00000	0.50000
O1	4a	0.23110	0.25000	0.11130
O12	4a	0.76890	0.75000	0.88870
O2	8b	0.15405	0.05340	0.70130
O22	8b	0.83595	0.55340	0.29870

General position:
x, y, z not restricted
by symmetry!

Magnetic space group: **P_bmn2₁** (31.129)

in non-standard setting.

to transform to conventional setting :

$(-b, a, c; 3/8, 1/4, 0)$

WP	+ (1' 1/2 0 0)
8b	$(x, y, z m_x, m_y, m_z), (-x+1/4, -y, z+1/2 -m_x, -m_y, m_z),$ $(x, -y+1/2, z -m_x, m_y, -m_z), (-x+1/4, y+1/2, z+1/2 m_x, -m_y, -m_z)$
4a	$(x, 1/4, z 0, m_y, 0), (-x+1/4, 3/4, z+1/2 0, -m_y, 0)$

Magnetic moments of the asymmetric unit (μ_B):

Mn1 **3.87** ≈ 0.0 ≈ 0.0

Split independent
positions in the lower
symmetry

zero values are not
symmetry "protected"

Magnetic Point Group: **mm21'**

MTENSOR: Symmetry-adapted form of crystal tensors properties of magnetic crystals

Magnetic Symmetry and Applications

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MTENSOR: Tensor calculation for Magnetic Point Groups

For the symmetry-adapted form of non-magnetic crystal tensors see TENSOR

Tensor calculation for Magnetic Point Groups

MTENSOR provides the symmetry-adapted form of tensor properties for any magnetic point (or space) group. On the one hand, a point or space group must be selected. On the other hand, a tensor must be defined by the user or selected from the lists of known equilibrium, optical, nonlinear optical susceptibility and transport tensors, gathered from scientific literature. If a magnetic point or space group is defined and a known tensor is selected from the lists the program will obtain the required tensor from an internal database; otherwise, the tensor is calculated live. Live calculation of tensors may take too much time and even exceed the time limit, giving an empty result, if high-rank tensors, and/or a lot of symmetry elements are introduced.

Tutorial of MTENSOR: [download](#)

Further information can be found [here](#)

If you are using this program in the preparation of an article, please cite this reference:

Gallego *et al.* "Automatic calculation of symmetry-adapted tensors in magnetic and non-magnetic materials: a new tool of the Bilbao Crystallographic Server" *Acta Cryst. A* (2019) **75**, 438-447.

Please, enter a magnetic point group or a magnetic space group:

Magnetic Point or Space Group number:

Please, choose a tensor by one of these ways:

Choose a tensor from the lists

- Show symmetry-adapted tensors for all the magnetic point groups in standard setting
(this overrides previous choices)

EQUILIBRIUM TENSORS

OPTICAL TENSORS

NONLINEAR OPTICAL SUSCEPTIBILITY TENSORS

TRANSPORT TENSORS

Build your own tensor

- Introduce Jahn's symbol without superscripts. Examples: (1) $[[V_2][V_2]]$, (2) $a\{V_2\}^*$, (3) $(V_2[V_2])^*$

Detailed information in. Gallego et al., *Acta Cryst. A* (2019) 75, 438-447.
and tutorial: Tutorial_magnetic_section_BCS_1.pdf

MTENSOR

Magnetolectric tensor:

Group 6/m' (#23.4.85)

α^T_{ij}	j			
	1	2	3	
i	1	α^T_{11}	α^T_{12}	0
	2	$-\alpha^T_{12}$	α^T_{11}	0
	3	0	0	α^T_{33}

Number of independent coefficients: 3

Group 622 (#24.1.87)

α^T_{ij}	j			
	1	2	3	
i	1	α^T_{11}	0	0
	2	0	α^T_{11}	0
	3	0	0	α^T_{33}

Number of independent coefficients: 2

Group 62'2' (#24.4.90)

α^T_{ij}	j			
	1	2	3	
i	1	0	α^T_{12}	0
	2	$-\alpha^T_{12}$	0	0
	3	0	0	0

Number of independent coefficients: 1

Group 6mm (#25.1.91)

α^T_{ij}	j			
	1	2	3	
i	1	0	α^T_{12}	0
	2	$-\alpha^T_{12}$	0	0
	3	0	0	0

Number of independent coefficients: 1

Group 6m'm' (#25.4.94)

α^T_{ij}	j			
	1	2	3	
i	1	α^T_{11}	0	0
	2	0	α^T_{11}	0
	3	0	0	α^T_{33}

Number of independent coefficients: 2

Group -6'm'2 (#26.3.97)

α^T_{ij}	j			
	1	2	3	
i	1	α^T_{11}	0	0
	2	0	α^T_{11}	0
	3	0	0	α^T_{33}

Number of independent coefficients: 2

Group -6'm'2' (#26.4.98)

α^T_{ij}	j			
	1	2	3	
i	1	0	α^T_{12}	0
	2	$-\alpha^T_{12}$	0	0
	3	0	0	0

Number of independent coefficients: 1

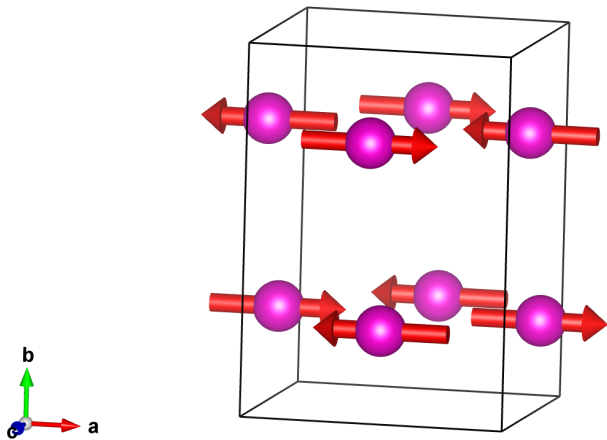
Group 6/m'mm (#27.3.102)

α^T_{ij}	j			
	1	2	3	
i	1	0	α^T_{12}	0
	2	$-\alpha^T_{12}$	0	0
	3	0	0	0

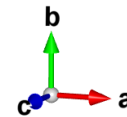
Number of independent coefficients: 1

Consequences of symmetry

EuZrO₃: [magndata #0.146 & 0.147](#)



Pnm'a



Pn'm'a'

Table of tensor components

α^T_{ij}	j		
	1	2	3
i	1	0	α^T_{13}
	2	0	0
	3	α^T_{31}	0

Number of independent coefficients: 2

Information about the selected tensor

- 2nd rank Magnetoelectric tensor α^T_{ij} (inverse effect)
- Axial tensor which inverts under time-reversal symmetry operation
- Defining equation: $\mathbf{P}_i = \alpha^T_{ij} \mathbf{H}_j$
- Relates Magnetic field \mathbf{H} with Polarization \mathbf{P}
- Intrinsic symmetry symbol: aeV^2

Output of MTENSOR

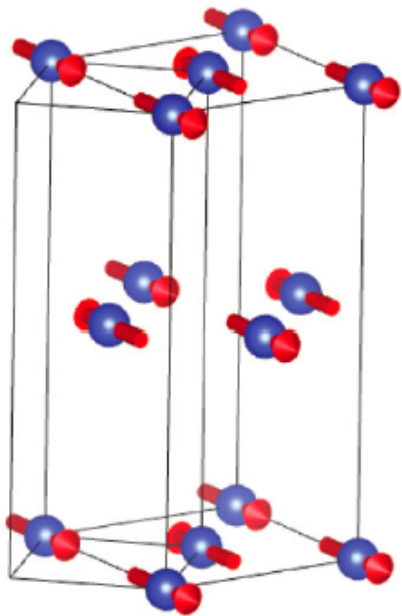
Table of tensor components

α^T_{ij}	j		
	1	2	3
i	1	α^T_{11}	0
	2	0	α^T_{22}
	3	0	0

Number of independent coefficients: 3

The same spin arrangement can produce different MSGs (and different ferroic properties) (*The non magnetic atoms are also important for the magnetic symmetry!*)

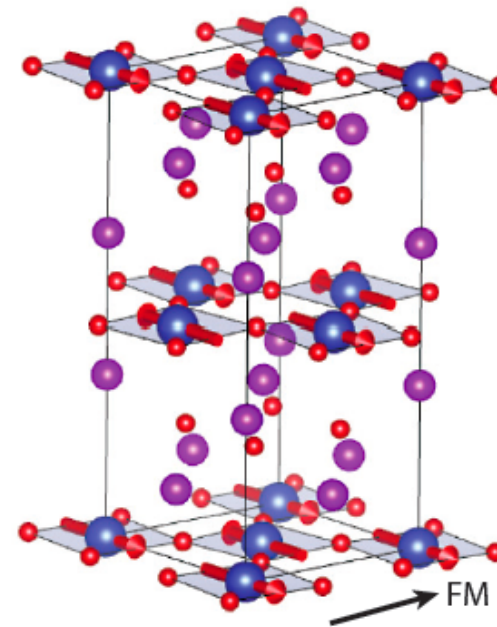
Pr₂CuO₄
I4/mmm, k=(1/2,1/2,0)



$C_{4v}ccm$
 (c, a - b, a + b; 1/4, 3/4, 1/4)

Point group: mmm1'

Gd₂CuO₄
Cmce, k=(0,0,0)



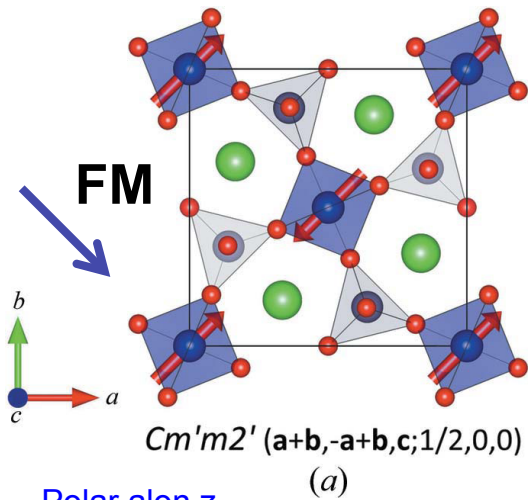
$Cm'ca'$
 (c, b, -a, 0, 0, 0)

Point group: m' mm'
 (weak ferromagnetism)

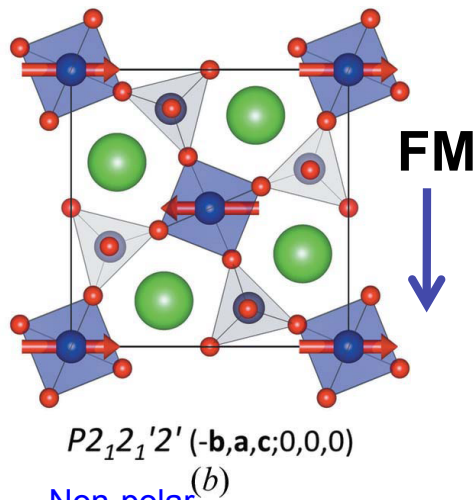
Consequences of symmetry

Ba₂CoGe₂O₇

parent SG: P-42₁m

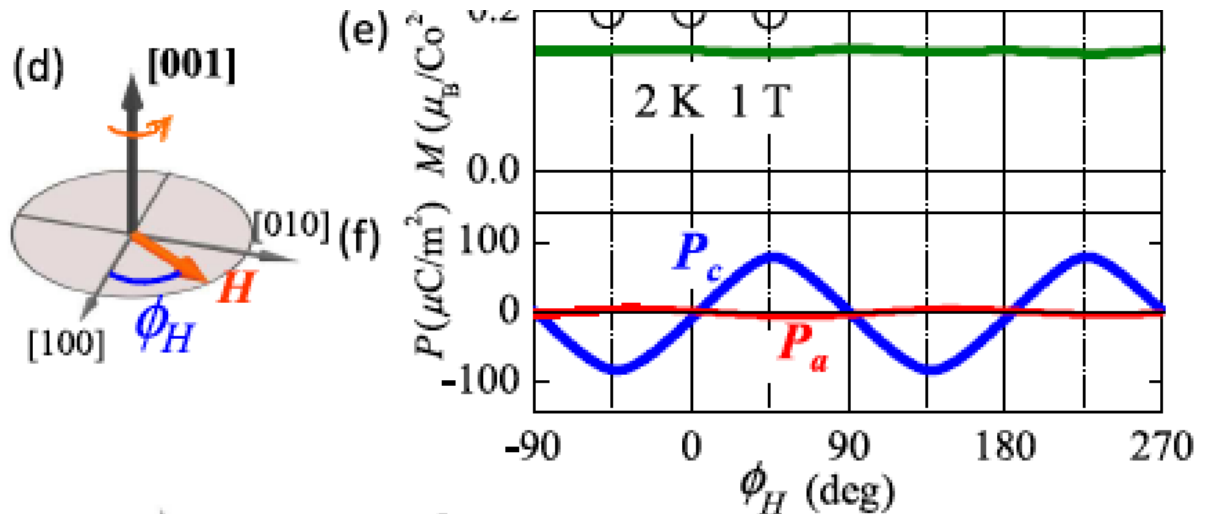


Polar along z
FM canting along $(-1, 1, 0)$



Non-polar
FM canting along $(0, 1, 0)$

Murakawa et al. PRL (2010):



Tutorial 1 :

Tutorial_magnetic_section_BCS_1

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